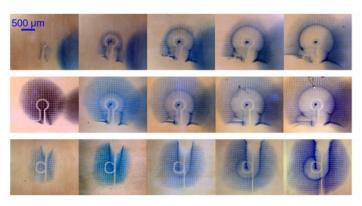


Dispersion Relation and Pulse Shaping of Femtosecond Laser-Driven Transient Pulsed Electromagnetic Fields



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Expert technical support and physical oversight:

C. Vlachos, M. Krupka, Ph. Korneev, G. Gatti, T. Pisarczyk, J. J. Santos

















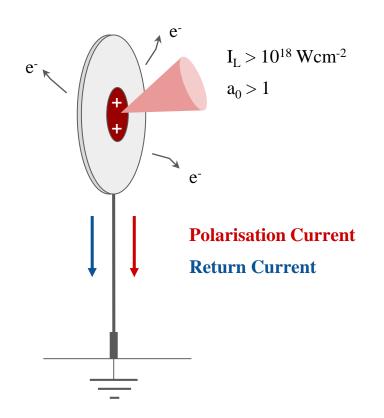
1. Introduction to laser-driven EM discharges

2. Our work on fast polarisation pulses driven by relativistic lasers

3. Our experiment on ELI-NP

Relativistic laser-solid interactions produce **electromagnetic pulses** (**EMP**) that propagate in "vacuum" and along the surface of the target





Hot electrons ejected by the laser polarise the target

[A. Brantov et al. PRE 2020]

Polarisation pulse propagates over target surface and EM pulse in free space

[K. Quinn et al. PRL 2009] [M. Ehret et al. PoP submitted] [S. Kar et al. Nat. Comm. 2016]

A return current is drawn when the polarisation wave reaches the ground

[F. Consoli et al. HPLSE 2020]

Antenna radiation is emitted when the return current passes along the target support

[J.-L. Dubois *et al.* PRE 2014] [A. Poyé *et al.* PRE 2015]

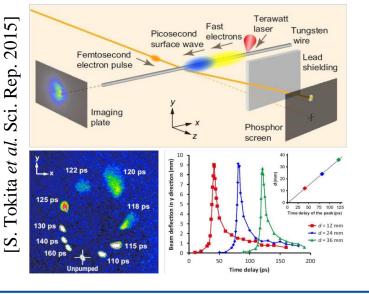
The fast polarisation pulse ($t \lesssim 50 \text{ ps}$)



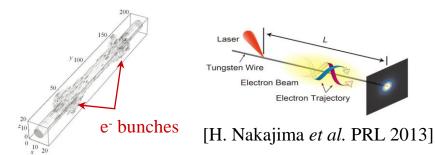
As hot e^{-s} cross target-vacuum boundary they excite an EM pulse that propagates in the skin-layer of the target at $v \sim c$

[K. Quinn et al. PRL 2009]

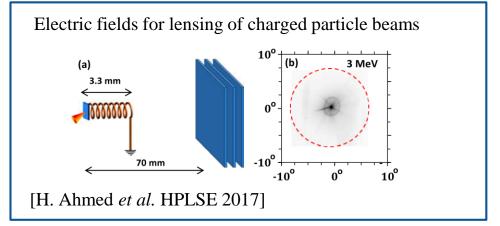
 $\tau \sim 10$ ps Sommerfeld-type EM surface wave Low dispersion over cm-distances



Relativistic e⁻s can be entrained within the surface wave



[A. Brantov *et al.* EPJ Web of Conferences 2018]

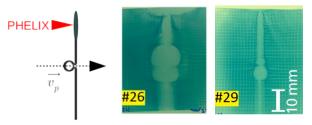


The Return Current ($t \ge 100 \text{ ps}$)

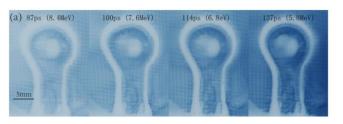


500 fs laser @ 5 x 10^{18} Wcm⁻² => ~ few kA / tesla in coil

30 fs laser @ 10^{18} Wcm⁻² => 20 T coil field

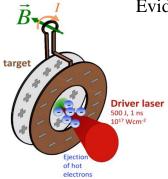


[M. Ehret *et al.* PoP submitted]

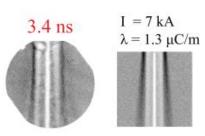


[W. Wang *et al.* PoP 2018]

A similar phenomenon is observed at lower intensity ($\lesssim 10^{17} \,\mathrm{Wcm^{-2}}$) and longer pulse duration (t_{las}~ ns), where the return current can produce quasi-static B-fields

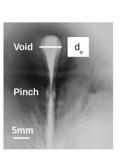


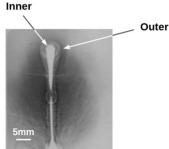
Evidence of multi-kA currents, wire plasma...



[M. Manuel et al. APL 2012]

...and long-lived GV/m E-fields





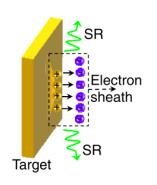
[J. J. Santos et al. PoP 2018]

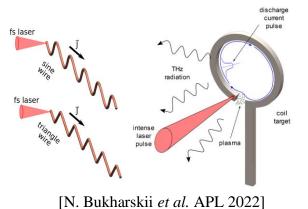
[P. Bradford et al. PPCF 2021]

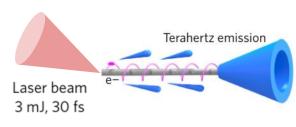
Discharge currents as secondary sources of radiation

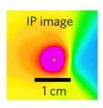


Bright THz sources







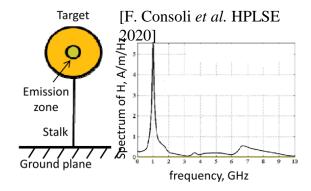


[Y. Tian et al. Nat. Phot. 2017]

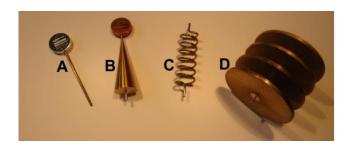
[G.-Q. Liao et al. PRX 2020]

[N. Bukharskii and Ph. Korneev ArXiV 2022]

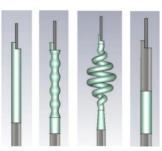
GHz EMP



[D. Minenna et al. PoP 2020]



[P. Bradford et al. HPLSE 2018]





1. Introduction to laser-driven EM discharges

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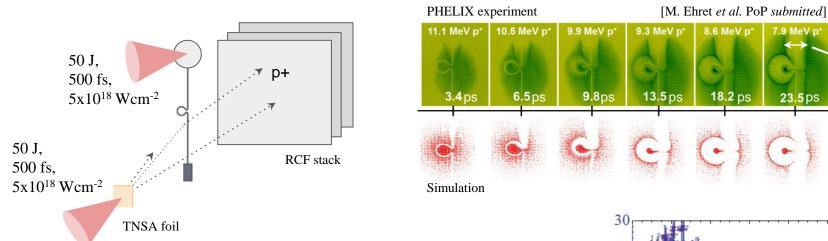
3. Our experiment on ELI-NP

Sub-c group velocity and Weibull-like "shape" of fast discharges measured with proton radiography

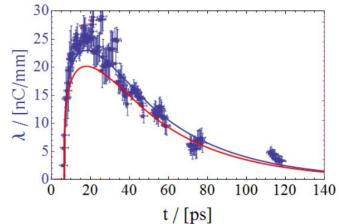


CAUSTIC

WIDTH

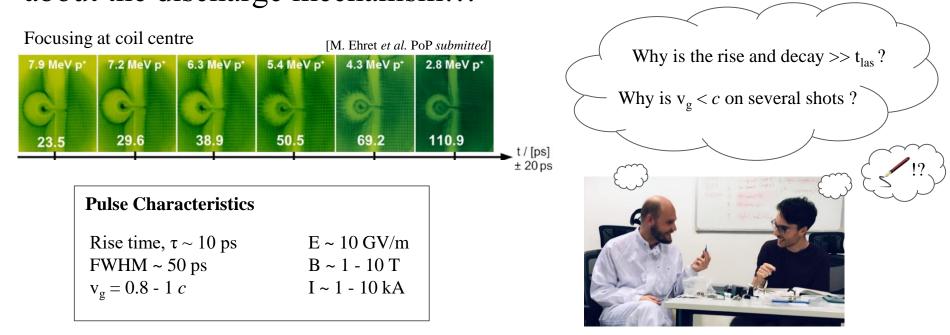


- 1. Analytical deflections dominated by radial E-field: Caustic width => Weibull-like shape
- 2. Electrodynamic p+ tracing simulation => amplitude, Q = 300 nC
- 3. Group velocity, $v_g \sim 0.8 1 c$



Proton lensing platform provokes interesting questions about the discharge mechanism...





... that are <u>important to answer</u> if we are to develop EM lenses or secondary sources of x-rays, GHz and THz radiation

"Sun-rays" and sub-c group velocities indicative of <u>plasma</u> formation on the wire surface?



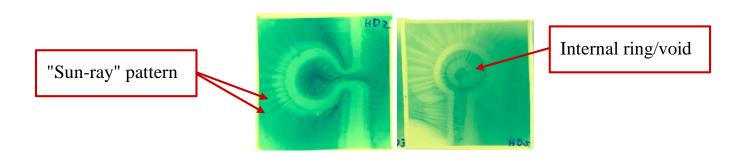
Conjecture 1

For a Sommerfeld wave propagating along a solid conductor, have $v_g \sim c$

Sub-c velocity \Rightarrow plasma formation on wire surface?

Conjecture 2

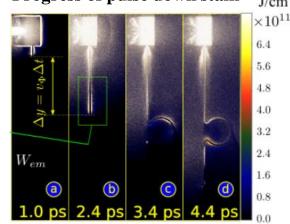
Potential modulations in p+ radiography \Rightarrow surface roughness can modify wave propagation speed and phase?



PIC sims reveal sub-c group velocity of discharge and possible plasma formation on wire

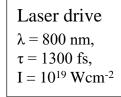






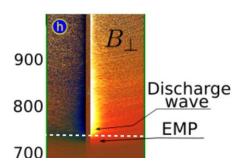
[M. Ehret et al. PoP submitted]

900 mm



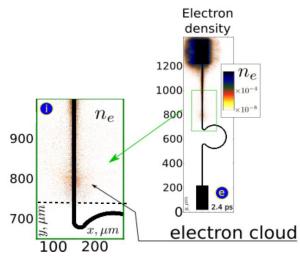
 $w_{\rm em} = B_{\perp}^2 / 2\mu_0 + \varepsilon_0 E_{\parallel}^2 / 2$ J/cm^2

Surface EM pulse propagates more slowly than the airborne EMP



B-field/current rise-time $\tau \sim 1$ ps is too fast to probe with p+ radiography

Plasma formation on wire and/or directly accelerated e⁻s?



ChoCoLaTII: Simplified Numerical Modelling of Target Charging



[A. Poyé et al. PRE 2018]

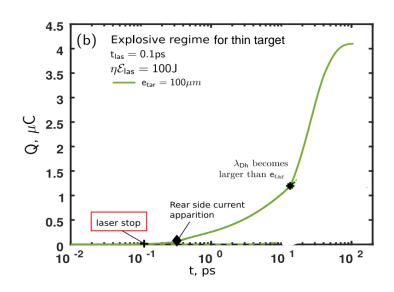
Target charging is governed by competition between the hot electron temperature T_e and target potential $\Delta\Phi$

average vel. of electrons that escape the barrier $\Delta\Phi$

$$I_h = e\Omega_b \left(\frac{\eta_\Phi N_h}{V_h}\right) \pi R_h^2 \int_{\Delta\Phi}^{\infty} f_h(\epsilon) v_h d\epsilon$$
 # density of electrons with energy $> \Delta\Phi$

Target potential $\Delta\Phi$ can be divided into thermal (from sheath) and electrostatic (from accumulated charge) components

ChoCoLaTII: At high intensity, the rise time of the charge profile can be $>> t_{las}$

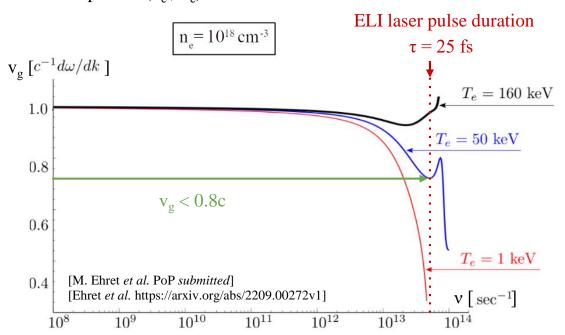


This is principally due to a <u>long hot e</u>-<u>collisional cooling time</u>

Analytic modelling of EM wave propagating in wire surface plasma suggests we may be able to tune pulse velocity and shape



Plasma Dispersion Relation: Frequency of polarisation pulse envelope related to plasma (n_a, T_a)



We can estimate the **frequency v** of the pulse by measuring:

plasma (n_e, T_e) group velocity, v_g

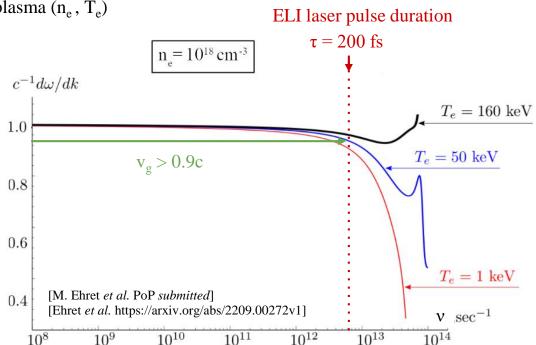
This allows us to discriminate between:

- Expanding surface charge (ChoCoLaTII - long rise-time)
- Sommerfeld wave induced by eejection from target $(\tau_{\text{pulse}} \simeq \text{laser pulse duration})$

Analytic modelling of EM wave propagating in wire surface plasma suggests we may be able to tune pulse velocity and shape



Plasma Dispersion Relation: Frequency of pulse envelope is related to plasma (n_e, T_e) ELI laser pulse duration



We can estimate the **frequency v** of the pulse by measuring:

plasma (n_e, T_e) group velocity, v_{σ}

This allows us to discriminate between:

- Expanding surface charge (ChoCoLaTII - long rise-time)
- Sommerfeld wave induced by e⁻ ejection from target $(\tau_{\text{pulse}} \simeq \text{laser pulse duration})$



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ELI-NP is needed to definitively answer the physics of target discharging



We have seen: p+ radiography is limited to ~1-10 ps resolution

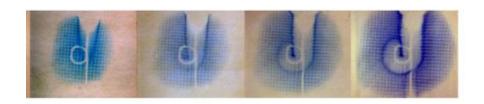
Need **sub-ps resolution** (of order t_{las}) to observe if the current rise time is $>> t_{las}$

Quantify impact of E_{las} , t_{las} and focal spot size on the shape (spatial and temporal) of the discharge pulse (ELI-NP ESSENTIAL)

Study the impact of plasma formation on \boldsymbol{v}_{g} of polarisation pulse

Study the effect of wire thickness, conductivity and roughness on $v_{\mathfrak{g}}$ and phase

Develop pulse trains for selective focusing



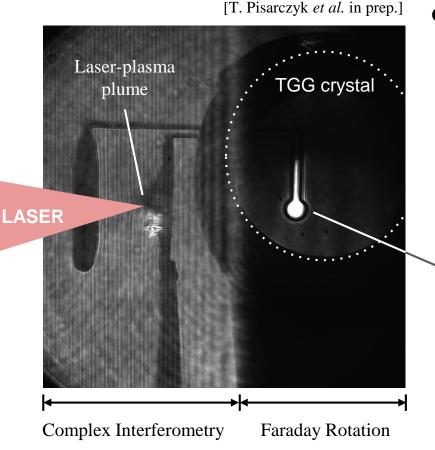
Focal spot scan
$$\Rightarrow$$
 $T_h \stackrel{?}{\Rightarrow} Q$

 E_{las} scan (fixed I_{las}) $\stackrel{!}{\Rightarrow}$ impact on Q and T_e via N_h

$$t_{las}$$
 (fixed I_{las} or fixed E_{las}) \Rightarrow effect on current and t_{pulse}

Using **optical probing** to diagnose wire currents and plasma conditions at sub-ps resolution





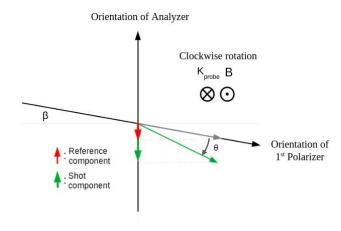
Coil Reference Image



Coil B-field induces polarisation rotation of probe beam

$$E_{probe} = 10 \text{ mJ}$$
 $\tau_{probe} \sim 40 \text{ fs}$
 $\lambda = 800 \text{ nm}$

Faraday Rotation Scheme



[T. Pisarczyk et al. J. Inst, 2019]

[C. Vlachos *et al*. in prep]

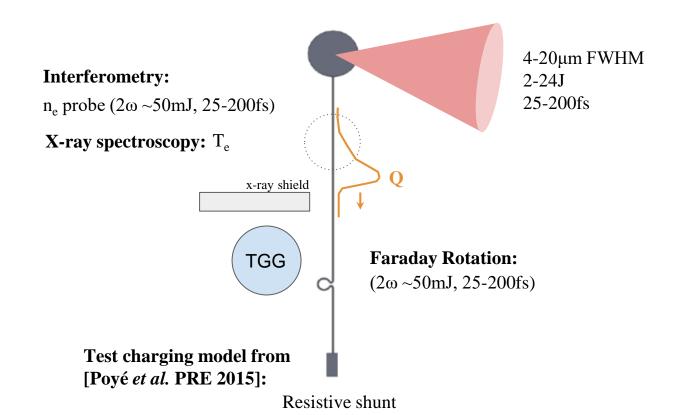
ELI-NP Proposal: First <u>fully-resolved</u> measurement of polarisation pulse and return current



Primary Diagnostics

Interferometry and \mathbf{x} -ray spectroscopy to infer plasma \mathbf{n}_{e} and \mathbf{T}_{e}

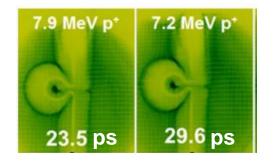
Faraday rotation measurement of polarisation/return currents



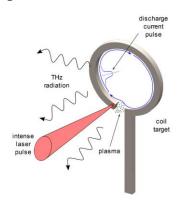
Conclusions



There remain important questions about the velocity, amplitude and shape of laser-driven currents...







[N. Bukharskii et al. APL 2022]

On ELI-NP:

Our experiment will measure the current evolution at >10x better temporal resolution than p+ radiography

We will examine how the discharge current is affected by the plasma dispersion relation of the wire

We want to use laser/target parameters to produce shaped kA current pulses and pulse trains